# A Study on Kicking Motion Strategy for a Legged Robot

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*Abstract*—For generating a translational motion of a robot in a horizontal plane or under microgravity, it is possible to utilize an impact force between the end-effector and a surface such as a wall and the ground. In this paper, we consider a legged robot consisting of four links that pushes a certain surface to obtain a translational momentum in a horizontal plane. The motion that maximizes the momentum is searched for, under the condition that the torque consumption during it is the same, by numerical optimization. The optimization results indicate that the motion with an impact force that is caused by hitting the surface is advantageous in increasing the translational momentum.

*Index Terms*—motion planning, differential evolution, impact force

## I. INTRODUCTION

Robots perform various tasks by acting on a target object or an environment such as the ground and walls with their end-effectors. For example, a legged robot can walk or jump by pressing its toes against the ground to achieve its locomotion. Those motions are highly dependent on the gravity, and large forces are necessary to support, accelerate or decelerate the robot in a vertical direction. For generating a vertical motion against gravity, it would be necessary to continuously push the ground for a certain duration, like a human walking or jumping [1-4].

In a horizontal plane or under microgravity, robots may achieve their locomotion by different motions that cannot be performed under gravity [5-7]. Due to recent space projects, researchers are increasingly interested in the locomotion of robots under microgravity [8-10]. It would be possible to generate an impact force between the endeffector and a surface to obtain a translational momentum of a robot. A novel way of locomotion could also be realized by choosing a robot mechanism that is specialized for microgravity. However, these possibilities have not been explored sufficiently.

Impact forces are usually reduced in motion control of robots [11,12], because they tend to cause mechanical damage. To mitigate the impact forces, several robot mechanisms have been proposed so far [13,14]. Although impact forces for the motions in a horizontal plane or under microgravity would be smaller than for a vertical motion under gravity, they can be reduced by attaching an elastic component on the end-effector that contacts with a surface.

In order to save the power consumption of a robot or increase the number of tasks that a robot can accomplish, we can choose an efficient way of motion according to the tasks [15-18]. Unlike humans and other living creatures, we can also use the robot mechanisms such as the joints that do not limit joint angles. The question then arises: what is the best way of kicking out a surface by a robot with 360 degrees rotation joints especially in a horizontal plane and under microgravity?

In this paper, we consider a four-link legged robot with 360 degrees rotation joints that pushes a surface in a horizontal plane. This corresponds to a humanoid robot pushing a wall horizontally, or a hopping robot under microgravity. We perform numerical simulations for the above mentioned two ways of motion, the motion of continuously pushing a surface and the motion of hitting it with an impact force. Under the condition that the torque consumption during operation is the same, the obtained momenta for the motions are compared to confirm the advantage of the motion with an impact force.

# II. DYNAMICAL MODEL OF LEGGED ROBOT

## A. Legged Robot

Fig. 1 shows a legged robot that was developed for experiments of jumping or pushing motion. The legged robot has four links and three joints, corresponding to the

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structure of the human leg. Body 1, Body 2, Body 3, and Body 4 represent the foot, lower leg, thigh, and torso, respectively. The legged robot kicks out a wall with the tip of its Body 1 (toe).



Figure 1. Legged robot

The model of contact force between the wall and the toe is important, because the translational momentum of the robot is obtained through the force. In the next subsection, we will present the equations of motion for the robot including the contact force.

## B. Equation of Motion

The inertial coordinate system and the body coordinate system for each Body *i* are introduced as shown in Fig. 2, where  $l_i$  is the length of the body and  $s_i$  is the distance from  $O_i$  to the center of gravity  $G_i$ . The point  $P_i$  of Body *i* is connected to the point  $O_{i+1}$  of Body i + 1 by a rotation joint. The equations of motion are given as follows.

$$\begin{bmatrix} \boldsymbol{M}_{1} & \boldsymbol{0} & \boldsymbol{0} & \boldsymbol{0} \\ \boldsymbol{0} & \boldsymbol{M}_{2} & \boldsymbol{0} & \boldsymbol{0} \\ \boldsymbol{0} & \boldsymbol{0} & \boldsymbol{M}_{3} & \boldsymbol{0} \\ \boldsymbol{0} & \boldsymbol{0} & \boldsymbol{0} & \boldsymbol{M}_{4} \end{bmatrix} \begin{bmatrix} \ddot{\boldsymbol{q}}_{1} \\ \ddot{\boldsymbol{q}}_{2} \\ \ddot{\boldsymbol{q}}_{3} \\ \ddot{\boldsymbol{q}}_{4} \end{bmatrix} = \begin{bmatrix} \boldsymbol{Q}_{1}^{v} \\ \boldsymbol{Q}_{2}^{v} \\ \boldsymbol{Q}_{3}^{v} \\ \boldsymbol{Q}_{4}^{v} \end{bmatrix} + \begin{bmatrix} \boldsymbol{Q}_{1}^{c} \\ \boldsymbol{Q}_{2}^{c} \\ \boldsymbol{Q}_{3}^{c} \\ \boldsymbol{Q}_{4}^{c} \end{bmatrix} + \begin{bmatrix} \boldsymbol{Q}_{1}^{e} \\ \boldsymbol{Q}_{2}^{e} \\ \boldsymbol{Q}_{3}^{e} \\ \boldsymbol{Q}_{4}^{e} \end{bmatrix}, \quad (1)$$

where  $M_i$  is an inertia matrix of Body *i* and is given as

$$\boldsymbol{M}_{i} = \begin{bmatrix} m_{i} & 0 & -m_{i}s_{i}\cos\phi_{i} \\ 0 & m_{i} & -m_{i}s_{i}\sin\phi_{i} \\ -m_{i}s_{i}\cos\phi_{i} & -m_{i}s_{i}\sin\phi_{i} & I_{i} + m_{i}s_{i}^{2} \end{bmatrix}.$$
 (2)

In the above equation,  $m_i$  and  $I_i$  are respectively the mass and the inertia of Body *i*. Its position and posture are denoted in a vector form as  $\boldsymbol{q}_i = [x_i \quad y_i \quad \phi_i]^T$ .

In the right-hand side of (1),  $\boldsymbol{Q}_{i}^{v}$  represents the inertia force, and is given as

$$\boldsymbol{Q}_{i}^{v} = \begin{bmatrix} -m_{i}s_{i}\dot{\varphi}_{i}^{2}\sin\varphi_{i} \\ m_{i}s_{i}\dot{\varphi}_{i}^{2}\cos\varphi_{i} \\ 0 \end{bmatrix}.$$
 (3)

 $Q_i^c$  represents the constraint force due to the coupling of the bodies at the joint, and is given as

$$\boldsymbol{Q}_{i}^{c} = -\left(\frac{\partial \boldsymbol{C}}{\partial \boldsymbol{q}_{i}}\right)^{T} \boldsymbol{\lambda}, \text{ where } \boldsymbol{C} = \begin{bmatrix} \boldsymbol{r}_{1}^{r} - \boldsymbol{R}_{2} \\ \boldsymbol{r}_{2}^{p} - \boldsymbol{R}_{3} \\ \boldsymbol{r}_{3}^{p} - \boldsymbol{R}_{4} \\ \boldsymbol{y}_{4} \\ \boldsymbol{\phi}_{4} \end{bmatrix} = \boldsymbol{0}. \quad (4)$$



Figure 2. Body coordinate system (left) and inertial coordinate system (right).

The first three components in C express the constraints at the rotation joints, where  $r_i^P$  is the vector from O of  $\Sigma_0$  to the point  $P_i$ , and  $R_i$  is a vector from O of  $\Sigma_0$  to  $O_i$  of  $\Sigma_i$ , as shown in Fig. 3. In addition, the *x*-direction motion and rotational motion of Body 4 are constrained for the robot in Fig. 1, to focus only on the motion perpendicular to the wall. The last two components in C correspond to those constraints.  $\lambda$  in (4) is the vector which represents the reaction forces due to the constraints.



Figure 3. Constraint at the rotation joint.

The external force  $Q_i^e$  is composed of the term from the actuator torques and the term from the contact force between the toe and the wall. Let  $\boldsymbol{\tau} = [\tau_1 \quad \tau_2 \quad \tau_3]^T$ , where  $\tau_i$  is an applied torque at Joint *i*. Then, the component from  $\boldsymbol{\tau}$  is expressed as  $H_i \boldsymbol{\tau}$ , where

$$H_{1} = \begin{bmatrix} 0 & 0 & 0 \\ 0 & 0 & 0 \\ 1 & 0 & 0 \end{bmatrix}, H_{2} = \begin{bmatrix} 0 & 0 & 0 \\ 0 & 0 & 0 \\ -1 & 1 & 0 \end{bmatrix}, H_{3} = \begin{bmatrix} 0 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & -1 & 1 \end{bmatrix}, H_{4} = \begin{bmatrix} 0 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & -1 \end{bmatrix}.$$
(5)

For i = 2,3,4, the external force is expressed as

$$\boldsymbol{Q}_i^e = \boldsymbol{H}_i \boldsymbol{\tau}. \tag{6}$$

For Body 1, we represent the force  $\boldsymbol{Q}_1^e$  as

$$\boldsymbol{Q}_{1}^{e} = \boldsymbol{H}_{1}\boldsymbol{\tau} + \boldsymbol{W}_{N}\boldsymbol{f}_{N} + \boldsymbol{W}_{T}\boldsymbol{f}_{T}, \qquad (7)$$

$$\boldsymbol{W}_N = \begin{bmatrix} 0 & 1 & 0 \end{bmatrix}^T, \boldsymbol{W}_T = \begin{bmatrix} 1 & 0 & 0 \end{bmatrix}^T,$$
 (8)

where  $f_N$  and  $f_T$  are the normal contact force and tangential force respectively. In this paper,  $f_N$  is represented by a virtual spring from the Hooke contact model [19], and  $f_T$  is expressed by a virtual damper.

$$f_N = \begin{cases} 0 & (g_N \ge 0) \\ -Kg_N & (g_N < 0) \end{cases},$$
(9)

$$f_T = \begin{cases} 0 & (g_N \ge 0) \\ -\mu f_N \text{sign}(\dot{g}_T) & (g_N < 0) \end{cases}, \quad (10)$$

where *K* is the stiffness coefficient and  $\mu$  is the friction coefficient.  $g_N$  and  $g_T$  are the normal and lateral direction distances between the toe of legged robot and the wall. The normal and tangential forces act only when the toe penetrates the wall.

### III. OPTIMIZATION OF KICKING MOTION

In this paper, we search for an optimal motion of the legged robot for kicking out the wall. This section summarizes the method of calculating an optimized motion numerically.

## A. Formulation of Optimization Problem

The torque  $\tau$  is assumed to be applied for  $t \in [0, t_f]$ , and to be kept at zero after the time  $t_f$ , as shown in Fig. 4. The legged robot gains translational momentum in the *y*direction from the wall only while the toe contacts with the wall. In other words, the greater the velocity obtained when taking off the wall, the greater the movement in unit time. Therefore, we consider an optimization problem that maximizes the velocity of the center of gravity in the *y*direction at the instant of taking off.

The optimization problem is formulated as follows.

$$\begin{cases} \text{Maximize} & \xi(\mathbf{X}) = \dot{G}_{y} \\ \text{w.r.t.} & \mathbf{X} = \begin{bmatrix} \mathbf{\tau}^{T} & t_{f} & \emptyset_{i}(0) \end{bmatrix} & (i = 1, 2, 3) & (11) \\ \text{s.t.} & \int_{0}^{t_{f}} (\tau_{1}^{2} + \tau_{2}^{2} + \tau_{3}^{2}) dt = \alpha \end{cases}$$

Denoting the position of gravity center of the robot along y axis as  $G_y$ , evaluation function  $\xi$  is given as its time derivative after taking off. The design variable **X** includes the torque  $\tau$ , the time  $t_f$  and the initial angles  $\emptyset_i(0)$  of Body 1 to 3. Those initial angles correspond to the initial posture of the robot. It should be noted that the torque consumption for  $t \in [0, t_f]$  is constrained to be the same amount of  $\alpha$ .

Since a time history of a torque is infinite-dimensional, we express it approximately by using the Fourier series.

$$\tau_i = a_{1,i} \sin\left(\frac{\pi t}{t_f}\right) + \dots + a_{h,i} \sin\left(\frac{h\pi t}{t_f}\right) \quad (12)$$

Then, the design variables  $\tau_i(t)$  are replaced by the coefficients  $a_{j,i}(i = 1, 2, 3, j = 1, \dots, h)$ . In addition, since it is difficult to satisfy the condition of the torques exactly, an acceptable range  $\beta$  is introduced.

$$\beta = ae^{-n} + b, \ \left(a = \frac{(\rho - 1)\varepsilon}{e^{-1} - e^{-L}}, \ b = \frac{e^{-1} - \rho e^{-L}}{e^{-1} - e^{-L}}\varepsilon\right),$$
(13)



Figure 5. Acceptable range  $\beta$  ( $\beta = \rho \varepsilon$  for = 1,  $\beta = \varepsilon$  for n = L).

Generations

where *n* and *L* are the number of generations and the maximum number of generations respectively, and will be described in the next subsection. The range  $\beta$  is narrowed down exponentially according to *n*, as shown in Fig. 5, where  $\rho(> 1)$  is a constant.

The evaluation function is also replaced by  $\bar{\xi}(X) = 1/\dot{G}_y$  to match the DE-based algorithm described in the next subsection. The optimization problem in (11) is rewritten as follows.

$$\begin{cases} \text{Minimize} & \xi(\mathbf{X}) = 1/G_{y} \\ \text{w.r.t.} & \mathbf{X} = [a_{j,i} \quad t_{f} \quad \emptyset_{i}(0)] \\ \text{s.t.} & J_{c} = \left| \int_{0}^{t_{f}} (\tau_{1}^{2} + \tau_{2}^{2} + \tau_{3}^{2}) dt - \alpha \right| \leq \beta \end{cases}$$
(14)

#### B. Algorithm Based on DE

In this paper, the optimization problem in (14) is solved with an algorithm based on Differential Evolution (DE) [20] that is described below.

(1) Generate each element of the *N* initial individuals  $X_k^{(1)}$  ( $k = 1, 2, \dots, N$ ) with a uniform random number of  $[a_{min}, a_{max}]$ .

(2) Repeat the following operations a) to d) with  $n = 1, 2, \dots, L$ , where *n* is called the number of generations.

a) If the condition for the torque comsumption is satisfied, compute  $\bar{\xi}(\boldsymbol{X}_{k}^{(n)})$  by numerical simulation. If not, set  $\bar{\xi}(\boldsymbol{X}_{k}^{(n)}) = \infty$ . Let the *m*-th individual  $\boldsymbol{X}_{m}^{(n)}$  for which

 $\bar{\xi}(\boldsymbol{X}_{m}^{(n)})$  is the smallest be the best individual  $\boldsymbol{X}_{best}^{(n)}$  $\left(\boldsymbol{X}_{best}^{(n)} = \boldsymbol{X}_{m}^{(n)}\right)$ . If  $\bar{\xi}(\boldsymbol{X}_{k}^{(n)}) = \infty$  for all  $k = 1, 2, \dots, N$ , the *m*-th individual for which  $J_{c}$  is the smallest is the best individual  $\boldsymbol{X}_{best}^{(n)}\left(\boldsymbol{X}_{best}^{(n)} = \boldsymbol{X}_{m}^{(n)}\right)$ .

b) Set the value of base vector  $X_{base}^{(n)}$  to be  $X_{best}^{(n)}$ . Also, two integers  $r_{1,k}$ ,  $r_{2,k}$  ( $r_{1,k} \neq r_{2,k} \neq m$ ) are randomly chosen from among  $k = 1, 2, \dots, N$  to generate the displacement vector  $M_k^{(n)}$  which is defined as follows.

$$\boldsymbol{M}_{k}^{(n)} = \boldsymbol{X}_{\text{base}}^{(n)} + F\left(\boldsymbol{X}_{r_{1,k}}^{(n)} - \boldsymbol{X}_{r_{2,k}}^{(n)}\right), \quad (15)$$

where F(> 0) is the scale factor.

c) Create a trial vector  $\boldsymbol{U}_{k}^{(n)}$ . Each element of  $\boldsymbol{U}_{k}^{(n)}$ ,  $U_{l,k}^{(n)}$   $(l = 1, 2, \dots, 3h + 4)$ , is defined as follows.

$$U_{l,k}^{(n)} = \begin{cases} M_{l,k}^{(n)} (R_{l,k} \le C \text{ or } l = l_{\text{rand}}) \\ X_{l,k}^{(n)} (\text{otherwise}) \end{cases}, \quad (16)$$

where  $l_{\text{rand}}$  is a randomly chosen integer from  $l, R_{l,k} \in [0 \ 1]$  is a uniform random number, and C(> 0) is a crossover rate.

d) Compute  $\bar{\xi}(\boldsymbol{U}_{k}^{(n)})$ . In this paper, if  $\bar{\xi}(\boldsymbol{X}_{k}^{(n)}) \neq \infty$ or  $\bar{\xi}(\boldsymbol{U}_{k}^{(n)}) \neq \infty$ , the surviving individual is selected according to (17). If  $\bar{\xi}(\boldsymbol{X}_{k}^{(n)}) = \infty$  and  $\bar{\xi}(\boldsymbol{U}_{k}^{(n)}) = \infty$ , it is selected according to (18).

$$\boldsymbol{X}_{k}^{(n+1)} = \begin{cases} \boldsymbol{U}_{k}^{(n)} \left( \bar{\xi}(\boldsymbol{U}_{k}^{(n)}) \leq \bar{\xi}(\boldsymbol{X}_{k}^{(n)}) \right) \\ \boldsymbol{X}_{k}^{(n)} \left( \bar{\xi}(\boldsymbol{U}_{k}^{(n)}) > \bar{\xi}(\boldsymbol{X}_{k}^{(n)}) \right) \end{cases}$$
(17)

$$\boldsymbol{X}_{k}^{(n+1)} = \begin{cases} \boldsymbol{U}_{k}^{(n)} \left( J_{c}(\boldsymbol{U}_{k}^{(n)}) \leq J_{c}(\boldsymbol{X}_{k}^{(n)}) \right) \\ \boldsymbol{X}_{k}^{(n)} \left( J_{c}(\boldsymbol{U}_{k}^{(n)}) > J_{c}(\boldsymbol{X}_{k}^{(n)}) \right) \end{cases}$$
(18)

#### IV. NUMERICAL RESULTS

#### A. Parameter Setting

In this section, the design variable X that minimizes  $\bar{\xi}(\mathbf{X})$  will be obtained by the above mentioned algorithm with the following parameters; The lengths and weights of the bodies are chosen as  $l_1 = 0.07$  (m),  $l_2 = l_3 =$ 0.30 (m),  $l_4 = 0.20$  (m),  $m_1 = 0.40$  (kg),  $m_2 = m_3 =$ 1.58 (kg), and  $m_4 = 1.42$  (kg). The inertia moment of each body is calculated under the assumption that each body is a uniform bar. The coefficients K in (9) and  $\mu$  in (10) are set to be  $K = 170 \times 10^5$  (N/m) and  $\mu = 0.1$ . Considering the output of the actuators in the experimental devise,  $\alpha$  in (11) is set to  $\alpha = 0.12$ .  $\varepsilon$  and  $\rho$  in (13) are chosen as  $\varepsilon = 0.012$  and  $\rho = 10$ . The torques are represented by (12) with h = 10, and the range of random numbers  $[a_{\min}, a_{\max}]$  is set to [-2.5, 2.5]. The number of individuals is N = 200, the number of maximum generation is L = 200, the crossover rate is C = 0.9, and the scale factor is F = 0.6.

For numerical calculation, we make the following assumptions; a) At the initial time t = 0, the legged robot is stationary, and the point  $O_1$  of Body 1 is in contact with the wall. b) The legged robot is assumed to slide on a smooth horizontal surface, and the friction between them is ignored. c) Joint friction is not considered. d) The initial angles  $\phi_i(0)$  of Body 1 to 3 are constrained as follows.

$$0 \le \phi_1(0) \le \frac{\pi}{2}, \ -\frac{\pi}{2} \le \phi_2(0) \le 0, \ 0 \le \phi_3(0) \le \frac{\pi}{2}$$
(19)

We consider two cases for numerical optimization, in order to compare two types of kicking motion: hitting the wall with an impact force and pushing the wall continuously. Although a motion such as repeatedly contacting the wall can be considered, the numerical optimization excludes it by evaluating  $G_y$  right after the toe first leaves the wall except at the initial time. To obtain the motion of pushing the wall continuously through numerical optimization, we introduce the following constraint.

$$\dot{\emptyset}_1 < 0 \text{ for } t > 0 \tag{20}$$

Numerical optimization was performed based on the algorithm in Sec. III in the following two cases: Case 1 without the constraint (20) and Case 2 with the constraint (20). In Case 2, we use the penalty method to find an optimized solution that satisfies the constraint (20), that is, if (20) is not satisfied, we make the evaluation function much worse.

#### B. Results and Discussion

The motions obtained by numerical optimization in Case 1 and Case 2 are shown in Fig. 6. The position and posture of the robot are drawn for both cases every 0.02 (s) from 0 (s) to 0.1 (s) and at 0.2 (s) in the figure, where the black dots indicate the position of the center of gravity. The torques for the motions are shown in Fig. 7, where the red, blue and green lines correspond to  $\tau_1$ ,  $\tau_2$ , and  $\tau_3$  respectively. The time histories of normal force  $f_N$  caused at the contact point are also shown in Fig. 8.

In Case 1, the legged robot makes almost one rotation of Body 1 counterclockwise around Joint 1 as shown in Fig. 6 (a), and then hits the wall around t = 0.087 (s) to obtain a translational momentum through a large impact force shown in Fig. 8 (a). The velocity of center of gravity  $\dot{G}_y$ , that is,  $\xi(\mathbf{X})$ , is 1.544 (m/s) immediately after the toe leaves the wall at t = 0.0874 (s), and the height of robot  $G_y$  reaches 0.63 (m) at t = 0.2 (s).





Figure 6. Position and posture obtained by numerical optimization.









In Case 2, the legged robot rotates Body 1 clockwise and pushes the wall continuously with the toe from the initial time, as shown in Fig. 6 (b). The contact force  $f_N$  is also caused continuously as shown in Fig. 8 (b) until the robot takes off at t = 0.0659 (s). The velocity of center of gravity  $\dot{G}_y$ , that is,  $\xi(\mathbf{X})$ , is 0.436 (m/s) after the take off, and the height of robot  $G_y$  reaches 0.51 (m) at t = 0.2 (s).

From these results, the velocities  $\dot{G}_y$  obtained from the same amount of torque consumption are largely different in Case 1 and Case 2; the velocity in Case 1 is about 3.5 times greater than that in Case 2. It would indicate that the motion of hitting the wall with an impact force in Case 1 is advantageous in generating a large translational momentum in a horizontal plane or under microgravity.

From the time histories of joint torques in Fig. 7, almost all the allowable torque consumption is utilized to rotate Body 1 by  $\tau_1$  in both cases. In Case 1, the torque  $\tau_1$  can generate a large rotational energy of Body 1, because the toe does not contact with the wall until the hitting and the rotation of Body 1 is easily accelerated. The energy accumulated until the hitting causes a large translational momentum through the elastic collision with the wall. On the other hand, in Case 2, the rotational acceleration of Body 1 is quite reduced due to the contact force from the wall. Even though the contact force is kept applied for a longer time, the total amount of translational momentum obtained from the contact force is smaller than in Case 1. Moreover, the vibratory behavior of the contact force in Case 2 is caused from the elasticity of the contact model.

The motion of hitting a surface with an impact force would be one promising strategy for kicking in a horizontal plane or under microgravity. Although a large impact force may cause mechanical damage to the robot, the peak of the force could be mitigated by utilizing an elastic component attached on the toe. Moreover, the impact force at each joint can be estimated by  $Q_i^c$  in the dynamic model presented in Sec. II. It should be noted that the postures of robot in both cases are almost stretched out, that is, close to singular configurations. The amount of contact force that can be obtained from the energy and joint torques of robot would depend on the posture. A more detailed analysis will be performed as a future work.

## V. CONCLUSION

In this paper, we searched for an optimal kicking motion of a legged robot in a horizontal plane by numerical optimization. The optimization results obtained by a DEbased algorithm showed that the translational momentum caused by hitting a surface with an impact force is much larger than the one by pushing it continuously, under the constraint that the torque consumption during the motion is constant. Accumulating the energy in an internal motion of the legged robot and transferring the energy to a translational momentum through a collision with a surface would be one of effective kicking motion strategies in a horizontal plane or under microgravity.

#### CONFLICT OF INTEREST

The authors declare no conflict of interest.

#### AUTHOR CONTRIBUTIONS

Masashi Nakamura conducted the research and wrote the paper; Takao Muromaki and Takateru Urakubo supervised and reviewed both research and paper; all authors had approved the final version.

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